Coupling-Dependent Data Flipping in Wireless Power and Data Transfer System

Hao Qiu, Takayasu Sakurai, and Makoto Takamiya Institute of Industrial Science The University of Tokyo, Tokyo, Japan Email: hqiu@iis.u-tokyo.ac.jp

Abstract—Load modulation is a common method to transfer data from the receiver (RX) side to the transmitter (TX) side in a magnetic resonance coupling wireless power transfer system. However, coupling-dependent data flipping (CDDF), that the data can be flipped depending on the coupling between TX and RX coils, is found for the first time using the conventional voltage monitoring across the TX coil. An analytic formula of the critical coupling coefficient (k_c) of the coils that determines CDDF is derived. To avoid CDDF, a TX input voltage monitoring method is proposed. Both the CDDF in the conventional method and the correct data transfer without CDDF in the proposed method are demonstrated in the measurement.

Keywords—Coupling-dependent data flipping, critical coupling coefficient, load modulation, magnetic resonance coupling, transmitter input voltage monitoring, wireless power transfer, wireless data transfer

I. INTRODUCTION

Wireless power transfer using magnetic resonance coupling has been applied in a wide range of subjects such as electric vehicles, mobile devices, and biomedical implants [1]. In a wireless power transfer system, when the coupling coefficient (k) between the TX and RX coils changes, both the coil-to-coil efficiency [2] and power losses in circuit components [3-4] are affected. In order to control and regulate the power in the RX coil, the establishment of a RX-to-TX communication (uplink) which is irrespective of k is required. As is shown in Fig. 1, in a wireless power and data transfer system, load modulation [5] has been widely used due to its simple architecture in which the change of the RX coil's loading condition is reflected in the TX coil and the voltage across the TX coil $(V_{\rm L})$ is monitored. However, most of the previous research [6-8] has focused on the case where k is a constant and few has discussed its possible effects.

II. CDDF AND THE SOLUTION

In this work, the effects of k on data transfer in a wireless power and data transfer system are discussed. As shown in Fig. 2(a), when k is smaller than a critical k (k_C), the data are demodulated depending on the amplitude of V_L . Here, data '1' is defined as large V_L and data '0' is defined as small V_L . When $k > k_C$, on the other hand, the relation between the data and the amplitude of V_L is flipped and incorrect data are demodulated, which we call coupling-dependent data flipping (CDDF). k_C is the critical point, at which V_L does not change its value during load modulation and thus data cannot be transferred from the



Fig. 1. Equivalent circuit model of a wireless power and data transfer system. The impedance matching networks (IMNs) in both TX and RX sides are not drawn for simplicity.



Fig. 2. Data backscatter based on (a) conventional $V_{\rm L}$ monitoring and (b) proposed $V_{\rm LC}$ monitoring.

RX side to the TX side. Thus, we propose to monitor the TX input voltage (V_{LC}) rather than V_L in Fig. 1. As shown in Fig. 2(b), where data '1' is defined as small V_{LC} and data '0' is defined as large V_{LC} , an uplink with a correct data transfer robust against CDDF is established.

III. CIRCUIT ANALYSIS OF CDDF

In this chapter, the origin of CDDF is analyzed. Fig. 1 shows the equivalent circuit model of a series-series topology,

where the inductors and compensation capacitors are connected in series in both TX and RX sides. The TX side, consisting of an inductor (L_{TX}) with an equivalent series resistance (r_{TX}) and a compensation capacitor (C_{TX}), is driven by a voltage source (V_S) with an internal resistance (R_S). The RX side is composed of an inductor (L_{RX}) with an equivalent series resistance (r_{RX}) and a compensation capacitor (C_{RX}). R_L is the load resistance. The currents flowing in the TX and RX sides are I_{TX} and I_{RX} , respectively.

Applying Kirchhoff's voltage law to the circuit, we obtain the following relations between the currents and voltages at the resonance frequency (f_0):

$$\begin{cases} \omega_{0} = 2\pi f_{0} = 1/\sqrt{L_{\text{TX}}C_{\text{TX}}} = 1/\sqrt{L_{\text{RX}}C_{\text{RX}}} \\ V_{\text{S}} = I_{\text{TX}}(R_{\text{S}} + r_{\text{TX}} + j\omega_{0}L_{\text{TX}} + 1/(j\omega_{0}C_{\text{TX}})) - j\omega_{0}MI_{\text{RX}} \\ j\omega_{0}MI_{\text{TX}} = I_{\text{RX}}(R_{\text{L}} + r_{\text{RX}} + j\omega_{0}L_{\text{RX}} + 1/(j\omega_{0}C_{\text{TX}})) \\ M = k\sqrt{L_{\text{TX}}L_{\text{RX}}} \end{cases}$$
(1)

Here, M is the mutual inductance between the TX and RX coils. j is the imaginary unit.

The ratio of $V_{\rm L}$ to $V_{\rm S}$ at f_0 can be described as:

$$\frac{V_{\rm L}}{V_{\rm S}} = \frac{I_{\rm TX}(j\omega_0 L_{\rm TX} + r_{\rm TX}) - I_{\rm RX} j\omega_0 M}{V_{\rm S}}
= \frac{(R_L + r_{\rm RX})(j\omega_0 L_{\rm TX} + r_{\rm TX}) + \omega_0^2 k^2 L_{\rm TX} L_{\rm RX}}{(R_{\rm S} + r_{\rm TX})(R_{\rm L} + r_{\rm RX}) + \omega_0^2 k^2 L_{\rm TX} L_{\rm RX}}.$$
(2)

Its amplitude is:

$$\left|\frac{V_{\rm L}}{V_{\rm S}}\right| = \frac{\sqrt{(R_{\rm L} + r_{\rm RX})^2 \omega_0^2 L_{\rm TX}^2 + \left[(R_{\rm L} + r_{\rm RX})r_{\rm TX} + \omega_0^2 k^2 L_{\rm TX} L_{\rm RX}\right]^2}}{(R_{\rm S} + r_{\rm TX})(R_{\rm L} + r_{\rm RX}) + \omega_0^2 k^2 L_{\rm TX} L_{\rm RX}}.$$
(3)

Based on the design parameters listed in Table I, a numerical calculation was performed. Fig. 3 shows k dependence of the calculated $|V_L/V_S|$ in the conventional V_L monitoring. k_C is defined when

$$|V_{\rm L} / V_{\rm S}|_{R_{\rm L}} = |V_{\rm L} / V_{\rm S}|_{R_{\rm L}=0\Omega}.$$
 (4)

It can be clearly seen that the data flips in the abnormal mode when $k > k_{\rm C}$ (~0.26). To understand the key parameters determining $k_{\rm C}$, we derived its analytic formula by assuming $r_{\rm TX}$ and $r_{\rm RX}$ as zero. That is:

$$k_{\rm C} = \sqrt{\frac{R_{\rm L}(\omega_0^2 L_{\rm TX}^2 - R_{\rm S}^2)}{2R_{\rm S}\omega_0^2 L_{\rm TX}L_{\rm RX}}}.$$
 (5)

Since ω_0 , L_{TX} , L_{RX} are determined by the coil design, we discuss the dependence of k_{C} on R_{L} and R_{S} . The R_{L} dependence of k_{C} is shown in Fig. 4(a), where R_{S} equals 50 Ω . Equation (5) gives the same positive R_{L} dependence as the numerical calculation based on (4), though with a deviation which comes from neglecting r_{TX} and r_{RX} . This deviation becomes smaller when R_{L} becomes larger. The R_{S} dependence of k_{C} is shown in Fig. 4(b), where R_{L} equals 10 Ω . When R_{S} becomes smaller than 5 Ω , k_{C} will become larger than 1 and CDDF cannot occur.

TABLE I. PARAMETERS OF TX AND RX COILS

	TX coil	RX coil
Coil diameter	300 mm	
Coil thickness	19 mm	14 mm
Wire diameter	1 mm	
Coil turns	10	
Resonance frequency (f ₀)	153 kHz	
Inductor	<i>L</i> _{τx} =66.2 μΗ	<i>L</i> _{RX} =73.8 μΗ
Equivalent series resistance (ESR)	<i>r</i> _{TX} =0.5 Ω	<i>r</i> _{RX} =1.0 Ω
Compensation capacitor	С _{тх} =16.4 nF	C _{RX} =14.7 nF
Quality factor	127	71



Fig. 3. k dependence of calculated $|V_{\rm L}/V_{\rm S}|$.

However, the impedance matching networks (IMNs) are usually inserted in both TX and RX sides to optimize the power transfer efficiency, and thus $R_{\rm S}$ seen from the TX coil's point of view equals $R_{\rm L}$ [9]. Thus, in the conventional $V_{\rm L}$ monitoring, CDDF cannot be avoided.

On the other hand, in the proposed V_{LC} monitoring, the amplitude of the ratio of V_{LC} to V_S at f_0 is described as:

$$\frac{\left|\frac{V_{\rm LC}}{V_{\rm S}}\right| = 1 - \frac{I_{\rm TX} r_{\rm TX}}{V_{\rm S}} = 1 - \frac{(R_L + r_{\rm RX}) R_{\rm S}}{(R_{\rm S} + r_{\rm TX})(R_{\rm L} + r_{\rm RX}) + \omega_0^2 k^2 L_{\rm TX} L_{\rm RX}}.$$
(6)

Fig. 5 shows k dependence of the calculated $|V_{LC}/V_S|$ in the proposed V_{LC} monitoring. It is seen that, irrespective of k,

$$|V_{\rm LC} / V_{\rm S}|_{R_{\rm L}} < |V_{\rm LC} / V_{\rm S}|_{R_{\rm L}=0\Omega}.$$
 (7)

That means CDDF never happens with V_{LC} monitoring and an uplink with a correct data transfer can be established.

Before discussing the measurement results, we give more explanations of why $k_{\rm C}$ exists in the conventional $V_{\rm L}$ monitoring and why CDDF never happens in the proposed $V_{\rm LC}$ monitoring. Fig. 6 gives the simplified circuit model of Fig. 1, assuming $r_{\rm TX}$ and $r_{\rm RX}$ to be zero. The equivalent resistance (r)



Fig. 4. (a) $R_{\rm L}$ and (b) $R_{\rm S}$ dependence of calculated $k_{\rm C}$.



Fig. 5. k dependence of calculated $|V_{LC}/V_S|$.

of the RX side seen from the TX side at f_0 can be described as [10]:

$$r = \omega_0^2 k^2 L_{\rm TX} L_{\rm RX} / R_{\rm L}.$$
 (8)



Fig. 6. Simplified circuit model of Fig. 1, in which r_{TX} and r_{RX} are assumed to be zero. *r* is the equivalent resistance of the RX side seen from the TX side.

 $|V_{\rm L}/V_{\rm S}|$ at f_0 is

$$\left|\frac{V_{\rm L}}{V_{\rm S}}\right| = \frac{\sqrt{r^2 + (\omega_0 L_{\rm TX})^2}}{r + R_{\rm S}}.$$
(9)

 $|V_{\rm L}/V_{\rm S}|$ equals 1 when $R_{\rm L}$ is shorted. When $R_{\rm L}$ is loaded, since *r* depends on *k*, $k_{\rm C}$ can exist to satisfy that

$$\sqrt{r^2 + \omega_0^2 L_{\text{TX}}^2} = r + R_{\text{S}}.$$
 (10)

At that point, $|V_L/V_S|$ also equals 1, which makes (4) hold.

On the other hand, $|V_{LC}/V_S|$ at f_0 is

$$\left|\frac{V_{\rm LC}}{V_{\rm S}}\right| = \frac{r}{r + R_{\rm S}}.\tag{11}$$

It can be easily seen that (7) holds based on (8) and (11).

IV. EXPERIMENTAL RESULTS

Fig. 7 shows the photo of the fabricated coils based on the parameters in Table I. A signal generator with $R_{\rm S}$ =50 Ω was used to generate a sinusoidal wave signal at 153kHz. *k* was measured according to the method described in [11] and efficiency (η) is defined as the ratio of the power directly transferred to $R_{\rm L}$ to the power extracted from the power source [10]. Fig. 8 gives the dependence of measured *k* and η on the distance (*d*) between the TX and RX coils.

Fig. 9 shows k dependence of measured $|V_L/V_S|$. k_C exists, as predicted in Fig.3. Compared with the calculation results, measured $|V_L/V_S|$ matches well when SW is Off but shows a deviation when SW is On. It can be ascribed to the unavoidable fact that the f_0 of the TX and RX coils can mismatch and the reflected reactance from the RX side will then affect the TX side. The effect is greater when SW is On. This is also the reason why the measured k_C (~0.43) in Fig. 9 does not match the calculated k_C (~0.26) in Fig. 3.

On the other hand, Fig. 10 shows the k dependence of measured $|V_{LC}/V_S|$ we proposed. As predicted in Fig. 5, (7) holds irrespective of k and CDDF never happens in the proposed V_{LC} monitoring. In addition, measured $|V_{LC}/V_S|$ in Fig. 10 matches calculation results in Fig. 5 well, which is different from the results of $|V_L/V_S|$ especially when SW is On. This is



Fig. 7. Photo of TX and RX coils.



Fig. 8. d dependence of measured k and η .

another advantage of the proposed $V_{\rm LC}$ monitoring that it supports an uplink robust against f_0 variations compared with the conventional $V_{\rm L}$ monitoring.

Here, more discussions are given. Series-series topology has been analyzed in this work, analyses in series-parallel, parallel-series, and parallel-parallel topologies will be included in our future work. In addition, the coil parameters were chosen mainly for EV applications. The universality of the phenomenon of CDDF will be discussed in several other application scenarios, such as mobile charging and biomedical applications.

V. CONCLUSIONS

To the best of our knowledge, the issue of CDDF was found for the first time in a wireless power and data transfer system. An analytic formula of $k_{\rm C}$ which shows a CDDF boundary was also derived. In order to solve CDDF, we proposed to monitor the transmitter input voltage ($V_{\rm LC}$). Circuit analysis and measurement were performed to demonstrate our conclusions.



Fig. 9. k dependence of measured $|V_{\rm L}/V_{\rm S}|$.



Fig. 10. k dependence of measured $|V_{LC}/V_S|$.

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